

Ages ago:

- 1st order bifurcation
- Intro to phase portraits of 2nd order systems.

Today:

- Continuation of phase portraits of 2nd order systems.
- Hartman-Grobman Thm
- Bendixon Thm
(Conditions for absence of periodic orbits of 2nd order systems) \rightarrow (just stands for 2nd order systems)

$$\dot{x} = Ax \xrightarrow{T} \dot{z} = \bar{A}z$$

$$(a) \quad \bar{A} = \begin{bmatrix} \lambda_1 & 0 \\ 0 & \lambda_2 \end{bmatrix}; \lambda_1 \neq \lambda_2, \lambda_1, \lambda_2 \in \mathbb{R}$$

$$(b) \quad \bar{A} = \begin{bmatrix} \lambda & 1 \\ 0 & \lambda \end{bmatrix}; \lambda \in \mathbb{R}$$

$$(c) \quad \lambda_{1,2} = \alpha \pm j\beta \rightarrow \bar{A} = \begin{bmatrix} \alpha & -\beta \\ \beta & \alpha \end{bmatrix}$$

(a) can show

$$z_2 = C z_1^{\frac{\lambda_2}{\lambda_1}}$$

$$\dot{z}_i = \lambda_i z_i \Rightarrow z_i(t) = e^{\lambda_i t} z_i(0) \Rightarrow \boxed{\text{if } z_i(0) = 0 \rightarrow z_i(t) = 0}$$

$i=1,2$

want $z_2 = z_2(z_1) \Rightarrow$ have to eliminate $t!$

$$\frac{z_2}{z_1} = \frac{z_2(0)}{z_1(0)} e^{(\lambda_2 - \lambda_1)t}$$

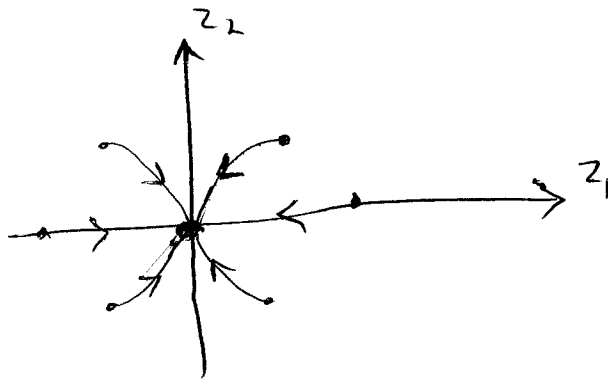
$$z_i(t) = e^{\lambda_i t} z_i(0) \rightarrow \frac{z_i(t)}{z_i(0)} = e^{\lambda_i t} \rightarrow t = \frac{1}{\lambda_1} \log \left(\frac{z_i(t)}{z_i(0)} \right)$$

... can show that

$$C = \frac{z_2(0)}{(z_1(0))^{\frac{\lambda_2}{\lambda_1}}}$$

Ex: $\lambda_1 < \lambda_2 < 0$

(Stable Node)

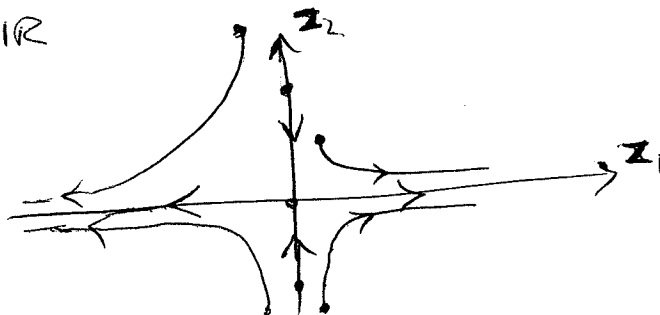


Another way of finding $z_2(z_1) \rightarrow$

$$\frac{\frac{dz_2}{dt}}{\frac{dz_1}{dt}} = \frac{\lambda_2 z_2}{\lambda_1 z_1} \rightarrow \boxed{\left(\frac{z_2}{z_1} \right)' = \frac{\lambda_2}{\lambda_1} \left(\frac{z_2}{z_1} \right)}$$

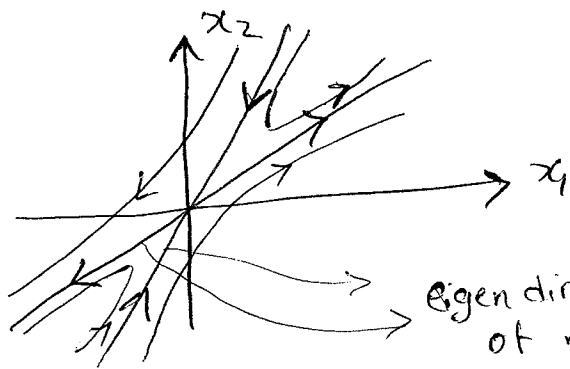
Ex: $\lambda_2 < 0 < \lambda_1 \in \mathbb{R}$

(Saddle Node)



$$Z = Tx \Rightarrow \begin{bmatrix} z_1 \\ z_2 \end{bmatrix} = \begin{bmatrix} T_{11} & T_{12} \\ T_{21} & T_{22} \end{bmatrix} \begin{bmatrix} x_1 \\ x_2 \end{bmatrix}$$

the n x - coordinates: (for saddle node)



→ these directions are eigenvectors!
[for z -coordinate: e_1 & e_2 are eigenvectors]

Aside:

Solutions of linear systems:

$$x(t) = e^{At} x(0)$$

$$z = Tx = T e^{At} T^{-1} z(0)$$

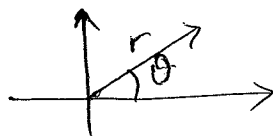
Fact: If A has a complete set of linearly independent eigenvectors, then we can choose: $T = V = [v_1 \dots v_n]$

$$\Rightarrow z(t) = e^{\Lambda t} z(0), \quad \Lambda = \text{diag}(\lambda_1, \dots, \lambda_n)$$

Case (C): $\bar{A} = \begin{bmatrix} \alpha & -\beta \\ \beta & \alpha \end{bmatrix} \rightarrow \lambda_{1,2} = \alpha \pm j\beta$

$$\begin{cases} \dot{z}_1 = \alpha z_1 - \beta z_2 \\ \dot{z}_2 = \alpha z_2 + \beta z_1 \end{cases}$$

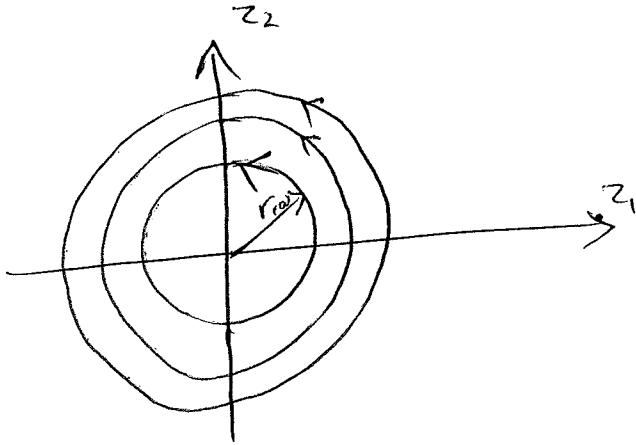
Polar Coordinates



$$\begin{cases} \dot{r} = \alpha r \\ \dot{\theta} = \beta \end{cases}$$

$$\begin{aligned} \rightarrow r(t) &= e^{\alpha t} r(0) \\ \rightarrow \theta(t) &= \beta t + \theta(0) \end{aligned}$$

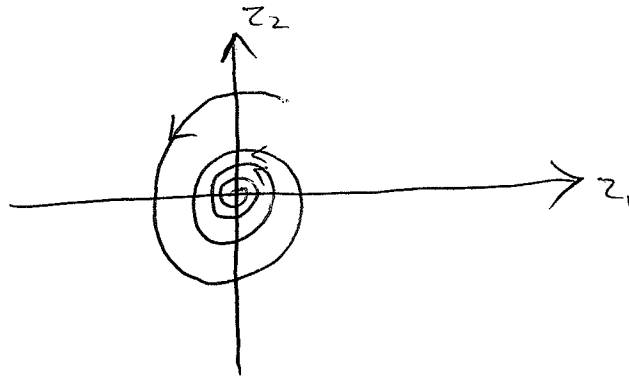
(The change of coordinate brings system to the even simpler one in comparison with z coordinates)



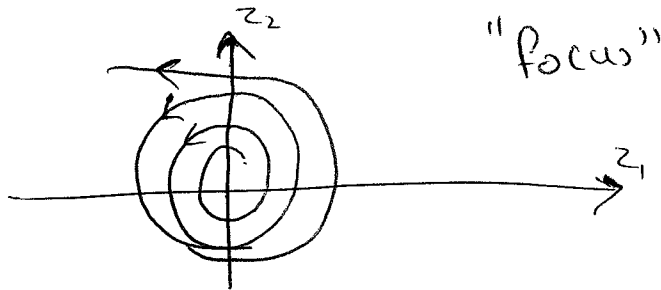
"Center"

$$\text{IF } \alpha = 0$$

$$\text{IF } \alpha < 0$$



$$\text{IF } \alpha > 0$$



"focus"

Case (b) :

$$\bar{A} = \begin{bmatrix} 0 & 1 \\ 0 & 0 \end{bmatrix} \Rightarrow$$

$$\ddot{y} = 0$$

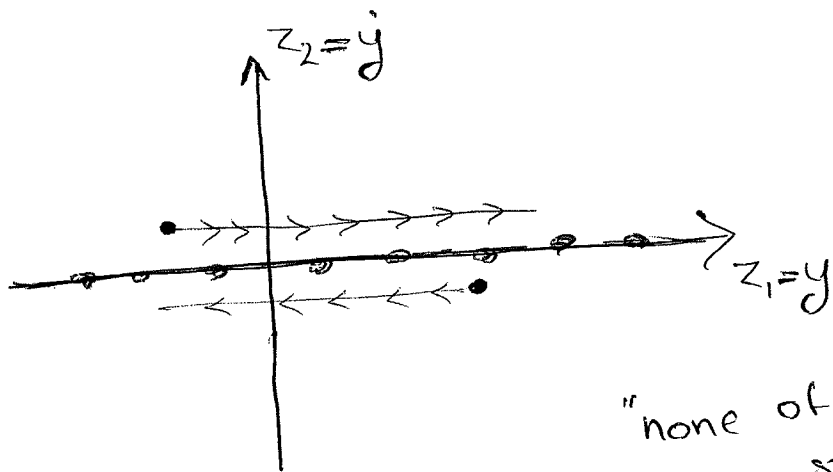
(moving mass on frictionless surface)

$$\rightarrow \begin{cases} \dot{z}_1 = z_2 & , & z_1 = y \\ \dot{z}_2 = 0 & , & z_2 = \dot{y} \end{cases}$$

$$z_2(t) = z_2(0)$$

$$z_1(t) = z_1(0) + z_2(0)t$$

$$\int \rightarrow \begin{cases} y(t) = y(0) + \dot{y}(0)t \\ \dot{y}(t) = \dot{y}(0) \end{cases}$$



"none of the eq. points are stable!"

Hartman - Grobman Thm

Notes:

If eq. point \bar{x} of $\dot{x} = f(x)$ is hyperbolic ($\bar{A} = \frac{\partial f}{\partial x} \Big|_{x=\bar{x}}$ doesn't have purely imaginary eigenvalues) then phase portraits of $\dot{x} = f(x)$ can be related to phase portrait of $\dot{\tilde{x}} = A\tilde{x}$

Thm: If \bar{x} is a hyperbolic eq. point of $\dot{x} = f(x)$ then there is a homeomorphism from a neighborhood of $\bar{x} \in \mathbb{R}^n$ that maps trajectories of $\dot{x} = f(x)$ to those of corresponding linearization.

Homeomorphism: Continuous map with continuous inverse!

Note: absence of e-values on $j\omega$ -axis of essence here!

(i.e. node, saddle, focus)

Ex:

$$\dot{x}_1 = -x_2 + ax_1(x_1^2 + x_2^2)$$

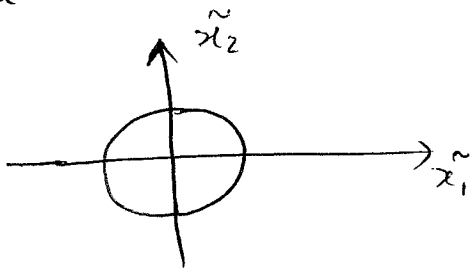
$$\dot{x}_2 = x_1 + ax_2(x_1^2 + x_2^2)$$

linearization around the origin: $\rightarrow \bar{x} = 0 = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$

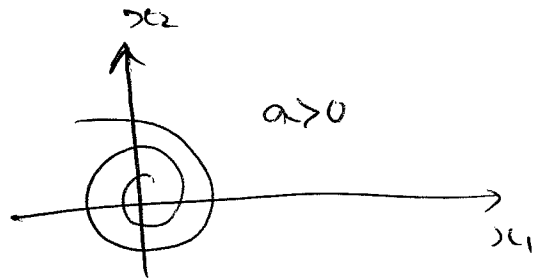
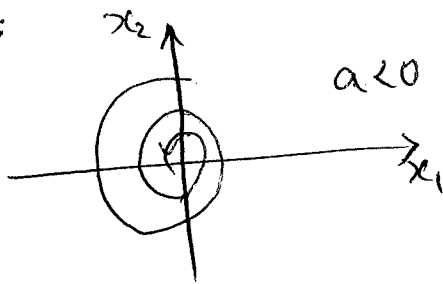
$$\rightarrow \begin{bmatrix} \tilde{x}_1 \\ \tilde{x}_2 \end{bmatrix} = \begin{bmatrix} 0 & -1 \\ 1 & 0 \end{bmatrix} \begin{bmatrix} \tilde{x}_1 \\ \tilde{x}_2 \end{bmatrix} \rightarrow \lambda_{1,2} = \pm j$$

(center)

linearization phase portrait:



Nonlinear:



polar coordinates: $\begin{cases} \dot{r} = ar^3 \\ \dot{\theta} = 1 \end{cases}$

If $a = 0 \rightarrow$ two systems are equal! but if $a \neq 0$

further calculations are required

So, the theorem works if we don't have any eigenvalues on imaginary axis.