

## Lecture 11 10/10

last time: modal decomp of LTI systems  
normal vs. non-normal matrices

today: recap of last time  
modal stability

E-value decomposition:

$$A \cdot v_i = \lambda_i \cdot v_i$$

↳ e-vector  
 ↳ e-value

A: diagonalizable

$$A = V \cdot \Lambda \cdot V^{-1}$$

if  $A \cdot A^T = A^T \cdot A$  (A: normal)

e-vectors of A are mutually orthogonal

$\Rightarrow$  A: unitarily diagonalizable

$$\Rightarrow A = V \cdot \Lambda \cdot V^*$$

↳ diagonal matrix  $\text{diag}\{\lambda_1, \dots, \lambda_n\}$

$$A \cdot f = V \cdot \Lambda \cdot V^* \cdot f = \left( \sum_{i=1}^n \lambda_i [v_i] \underbrace{[v_i^*]}_{\perp} \right) \cdot [f]$$

↳  $f \in \mathbb{R}^n$

$$v_i^* \cdot v_j = \delta_{ij} = \begin{cases} 1 & i=j \\ 0 & i \neq j \end{cases} \quad \rightarrow \text{sum of rank 1 matrices}$$

$$= \sum_{i=1}^n \lambda_i \underbrace{v_i^* \cdot f}_{\tau_i} [v_i]$$

↳ scalar

Thus for normal A: no intersection between different eigen directions

$$\text{For } \dot{x}(t) = Ax(t) \rightarrow \boxed{x(t)} = \sum \text{e}^{\lambda_i t} \left[ \begin{array}{c} v_i^* \\ \vdots \\ 1 \end{array} \right] \left[ \begin{array}{c} x_0 \\ \vdots \\ 1 \end{array} \right] \left[ \begin{array}{c} v_i \\ \vdots \\ 1 \end{array} \right] \quad \xrightarrow{\text{advancement in time}}$$

If  $A$  is nonnormal, ( $AA^T \neq A^TA$ ) then  $V^{-1} \neq V^*$

call it  $w^* \rightarrow V^{-1} = w^*$

i.e.  $A$  is NOT unitarily diagonalizable  
eigen-directions are NOT mutually orthogonal

$$Af = V \Lambda W^* = \sum_{i=1}^n \lambda_i [v_i] [w_i^*] \cdot [f]$$

$$= \sum_{i=1}^n \lambda_i w_i^* \cdot f \quad \begin{matrix} \text{not} \\ \text{orthogonal} \end{matrix}$$

$$x(t) = \sum_{i=1}^n e^{\lambda_i t} \cdot [w_i^*] [x_0] [v_i]$$

ex//  $A = \begin{bmatrix} -1 & 0 \\ k & -2 \end{bmatrix} \quad AA^T \neq A^TA$   
if  $k \neq 0$

$v_1 + v_2$  nearly parallel for large  $k$   
 $w_i^* A = \lambda_i w_i^*$  (left e-vector)

↳ we can always select them s.t.

$$w_i^* v_j = \delta_{ij} = \begin{cases} 1 & i=j \\ 0 & i \neq j \end{cases}$$

... Back to where we stopped last time

if there are complex e-values of  $A$  and  $A$  is diagonalizable  
then we can bring it to a block diagonal form:

$$\Lambda = \begin{bmatrix} \lambda_1 & & & \\ & \ddots & & \\ & & \lambda_r & \\ & & & \ddots & \lambda_m \end{bmatrix} \quad \{\lambda_1, \dots, \lambda_r\} : \text{real e-values of } A$$

$\Lambda_i \in \mathbb{R}^2 \times \mathbb{R}^2$  correspond to complex e-values

$$\lambda_i = z_i + j w_i$$

$$\Lambda_i = \begin{bmatrix} z_i & w_i \\ -w_i & z_i \end{bmatrix} \quad * \text{easy to show that e-value of } \Lambda_i \text{ are } z_i + j w_i$$

If  $\lambda_i \in \mathbb{C}$

$$c_i \cdot e^{\lambda_i t} \cdot v_i = z(t)$$

$\Rightarrow$  satisfies  $\dot{z}(t) = A \cdot z(t)$

since  $A \in \mathbb{R}^{n \times n}$ ,  $\operatorname{Re}(z(t))$  also satisfies  $\frac{d(\operatorname{Re}(z(t)))}{dt} = A \cdot \operatorname{Re}(z(t))$

$$\begin{aligned}\operatorname{Re}(z) &= [\operatorname{Re}(c_i) + j \operatorname{Im}(c_i)] \cdot e^{\lambda_i t} \cdot e^{j\omega_i t} \cdot [\operatorname{Re}(v_i) + j \operatorname{Im}(v_i)] \\ &\quad \cdot (\cos(\omega_i t) + j \sin(\omega_i t)) \\ &= \operatorname{Re}(z(t)) + \operatorname{Im}(z(t))\end{aligned}$$

Exercise

$$\text{Exercise: write } \operatorname{Re}(z(t)) = [\operatorname{Re}(v_i) \quad \operatorname{Im}(v_i)] \quad [ ] \quad \begin{bmatrix} \operatorname{Re}(c) \\ \operatorname{Im}(c) \end{bmatrix}$$

i.c. dependence

thus: "plane" spanned by  $\operatorname{Re}(v_i) + j\operatorname{Im}(v_i)$  invariant under system dynamics.

Question:

Can we diagonalize every matrix?

$$\text{ex// } A = \begin{bmatrix} \lambda & 1 \\ 0 & \lambda \end{bmatrix} \quad \lambda \in \mathbb{R}$$

E-values  $\lambda_1 = \lambda_2 = \lambda$

(algebraic multiplicity 2)

$$\lambda I - A = \begin{bmatrix} 0 & 1 \\ 0 & 0 \end{bmatrix}$$

single e-vector

$$v = \begin{bmatrix} * \\ 0 \end{bmatrix} \Rightarrow v = \begin{bmatrix} * & * \\ 0 & 0 \end{bmatrix}$$

Fact: Any matrix can be brought where  $J_i$  are given by:

$$\begin{bmatrix} J_1 & & & \\ & J_2 & \dots & \\ & & \ddots & J_m \end{bmatrix} \quad (*)$$

a.)  $\begin{bmatrix} \lambda_1 & & \\ & \ddots & \lambda_n \end{bmatrix}$  for  $\lambda_1, \dots, \lambda_n$  real + have linearly independent e-vectors

b.)  $\begin{bmatrix} w_i & w_i \\ -w_i & w_i \end{bmatrix}$  for complex  $\lambda_i$

c)  $\begin{bmatrix} \lambda & & 0 \\ 0 & \lambda & 0 \\ 0 & 0 & \lambda \end{bmatrix}$  for  $\lambda$ 's w/ linearly dependent e-vector  
 In several:  $\begin{bmatrix} \lambda & 0 & 0 \\ 0 & \lambda & 0 \\ 0 & 0 & \lambda \end{bmatrix}$

note: matrix of exp. of  $\star$ :  $\begin{bmatrix} e^{J_1 t} & & & \\ & e^{J_2 t} & & \\ & & \ddots & \\ & & & e^{J_m t} \end{bmatrix}$

ex//  $J = \begin{bmatrix} 0 & 1 \\ 0 & 0 \end{bmatrix} \Rightarrow e^{Jt} = \begin{bmatrix} 1 & t \\ 0 & 1 \end{bmatrix}$

ex//  $J = \begin{bmatrix} 0 & 1 & 0 \\ 0 & 0 & 1 \\ 0 & 0 & 0 \end{bmatrix} \Rightarrow e^{Jt} = \begin{bmatrix} 1 & t & \frac{1}{2}t^2 \\ 0 & 1 & t \\ 0 & 0 & 1 \end{bmatrix}$

ex// Jordan:  $\lambda_1 = -1$ ;  $\lambda_2 = -3$

$$\lambda_3 = 4 + 3j$$

$$\lambda_4 = 4 - 3j$$

$$\lambda_5 = \lambda_6 = \lambda_7 = 10$$

(+ 1 linearly independent e-vector)

$$\begin{bmatrix} J_1 & & \\ & J_2 & \\ & & J_3 \end{bmatrix} = \text{blkdiag}(J_1, J_2, J_3)$$

$$J_1 = \begin{bmatrix} -1 & 0 \\ 0 & -3 \end{bmatrix} \quad J_2 = \begin{bmatrix} 4 & 3 \\ -3 & 4 \end{bmatrix} \quad J_3 = \begin{bmatrix} 10 & 1 & 0 \\ 0 & 10 & 1 \\ 0 & 0 & 10 \end{bmatrix}$$

Summary: thus there is a change of coordinates to bring

$$\dot{x}(t) : A x(t) \text{ into } \dot{z}(t) = \bar{A} \cdot z(t) \text{ where } \bar{A} = \begin{bmatrix} J_1 & & \\ & \ddots & \\ & & J_m \end{bmatrix}$$

$$\Rightarrow z = \begin{bmatrix} z_1 \\ \vdots \\ z_n \end{bmatrix} \quad \dot{z}_i(t) = J_i \cdot z_i(t)$$

cannot do this w/ time-variant system