

Last time:

- Poincaré - Bendixon ThM
- Examples

Hopf Bifurcation

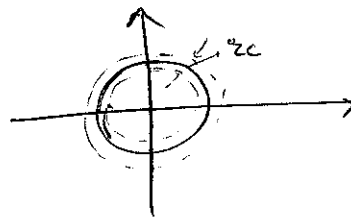
- Intro (Supercritical) 😊

Today

- Finish Hopf (Subcritical) 😞
- Non-dimensionalization
- Center manifold theory

Limit Cycle: an isolated closed trajectory V.S.

Eg.: Harmonic oscillator



there is another closed trajectory arbitrarily close to it!

(Not a limit cycle)

Supercritical :

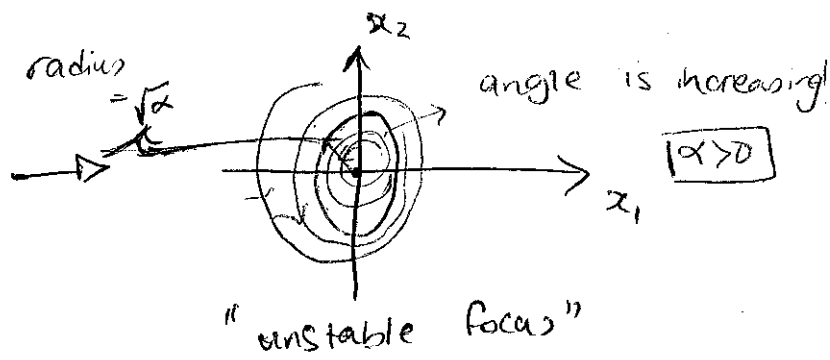
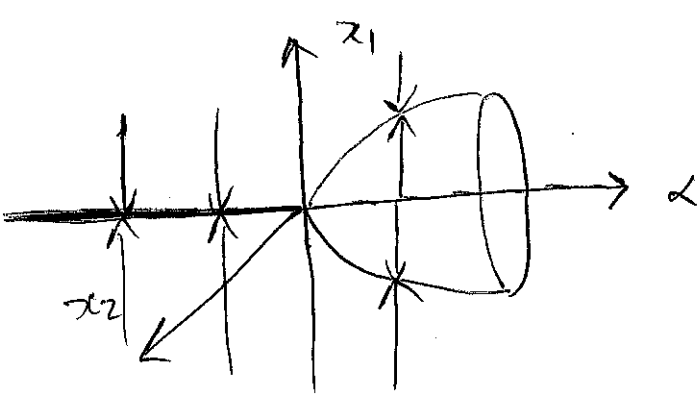
$$\begin{cases} \dot{r} = \alpha r - r^3 \\ \dot{\theta} = 1 \end{cases}$$

$$\rightarrow \dot{r} = f(r) \rightarrow f(r) = 0 \Leftrightarrow$$

it's a limit cycle

because  $\theta$  is spinning around!  $\rightarrow$  (this can never be an eq. point)

$$r = \begin{cases} 0 \\ \text{or} \\ \sqrt{\alpha}, \text{ if } \alpha > 0 \end{cases}$$

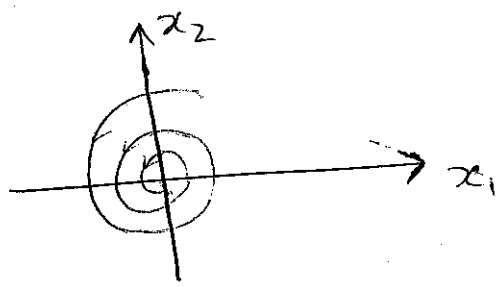


$$\dot{x}_1 = x_1 (\alpha - x_1^2 - x_2^2) - x_2$$

$$\dot{x}_2 = x_2 (\alpha - x_1^2 - x_2^2) + x_1$$

→ linearization  $\Rightarrow \begin{bmatrix} \dot{x}_1 \\ \dot{x}_2 \end{bmatrix} = \begin{bmatrix} \alpha & -1 \\ 1 & \alpha \end{bmatrix} \begin{bmatrix} x_1 \\ x_2 \end{bmatrix} \rightarrow \boxed{\lambda_{1,2} = \alpha \pm j}$

Fix  $\alpha < 0$



Sub critical Hopf:

$$\begin{cases} \dot{r} = \alpha r + r^3 - r^5 \\ \dot{\theta} = 1 \end{cases}$$

eq. points: for limit cycles  
(because angle is increasing,  
the origin is the only eq. point)

$$f(\bar{r}) = 0 \Rightarrow \bar{r} (\alpha + \bar{r}^2 - \bar{r}^4) = 0 \rightarrow \begin{cases} \bar{r} = 0 \text{ (eq. point)} \\ \bar{r}^4 - \bar{r}^2 - \alpha = 0 \end{cases} \rightarrow q = \bar{r}^2$$

$$\bar{r} \Rightarrow q^2 - q - \alpha = 0 \rightarrow q_{1,2} = \frac{1}{2} \{ 1 \pm \sqrt{1+4\alpha} \} \geq 0$$

$$= \bar{r}_{1,2}^2$$

Solutions exist iff:  $1+4\alpha \geq 0 \rightarrow \boxed{\alpha \geq -\frac{1}{4}}$

If  $\alpha = -\frac{1}{4} \rightarrow r_{12} = \frac{1}{2} \rightarrow \boxed{\bar{r} = \frac{1}{\sqrt{2}}}$

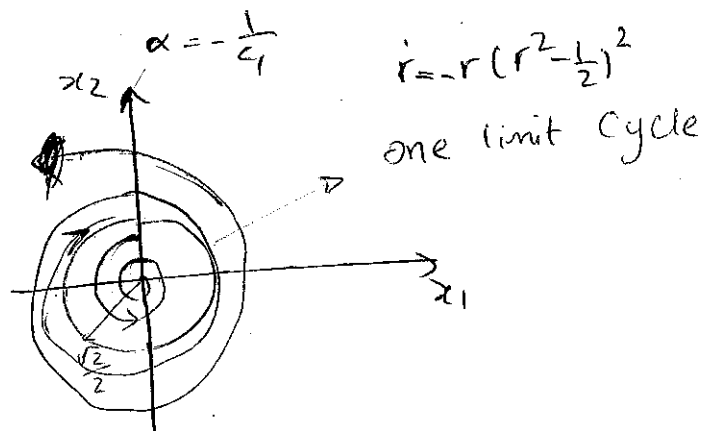
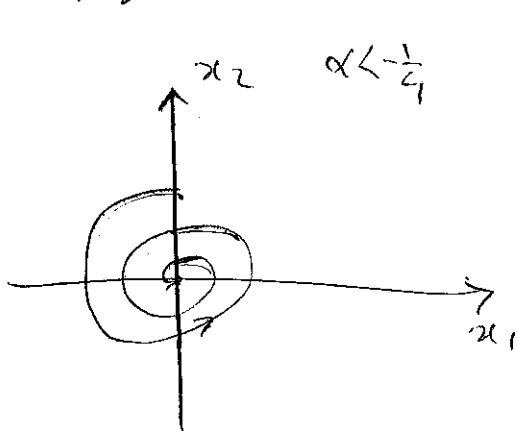
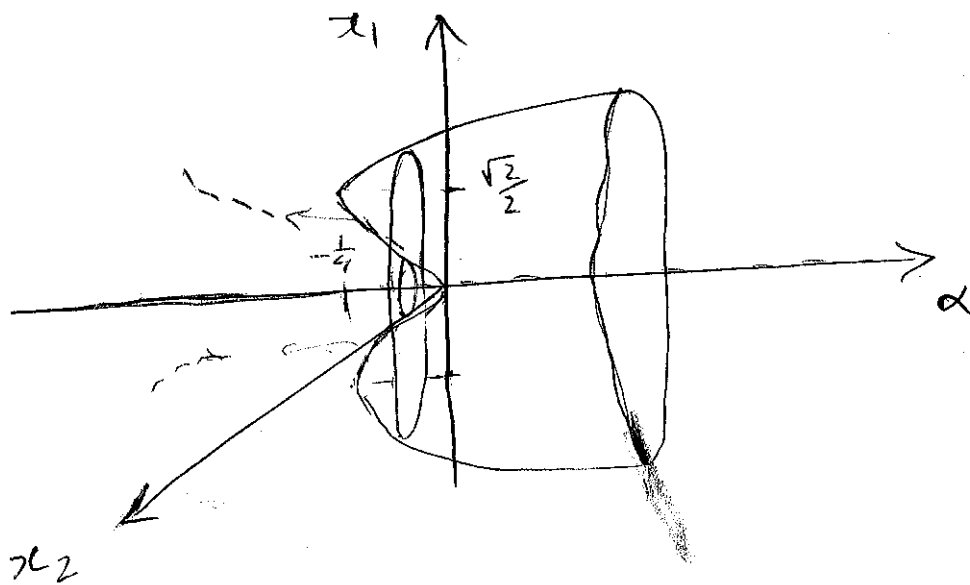
$\alpha < -\frac{1}{4} \rightarrow$  unique eq. point @ the origin!  
and no limit cycles

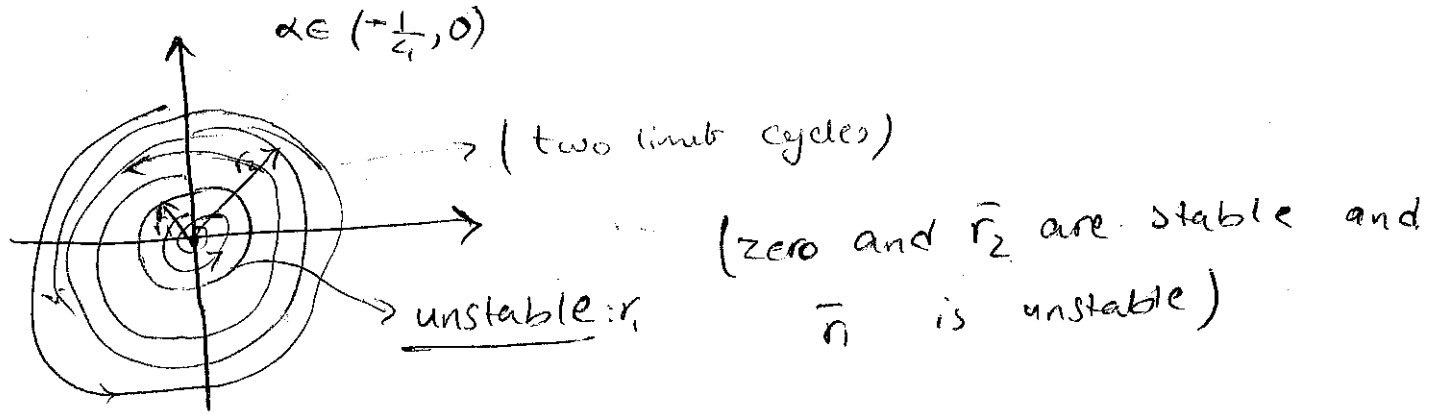
$\alpha = -\frac{1}{4} \rightarrow$  unique eq. point @ the origin!  
Limit cycle of radius  $\frac{1}{\sqrt{2}}$ !

$\alpha > -\frac{1}{4} \rightarrow$  unique eq. point @ the origin!  
two limit cycle  $\rightarrow \bar{r}_{12} = \sqrt{\frac{1}{2}(1 \pm \sqrt{1+4\alpha})}$   
cycle

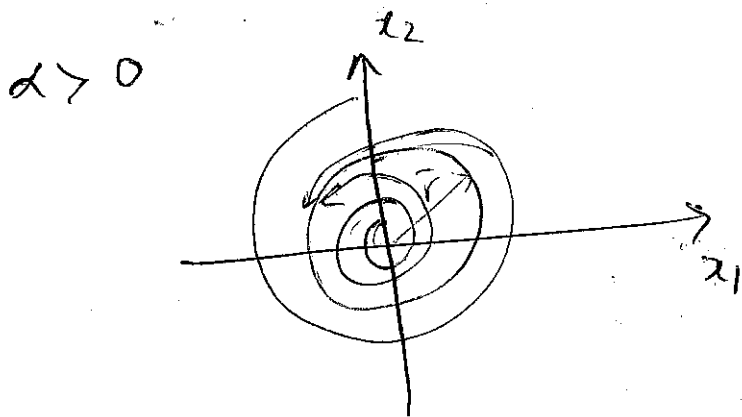
$\alpha = 0 \rightarrow$  unique eq. point @ the origin!  
Limit cycle of radius  $\bar{r} = 1$

$\alpha > 0 \rightarrow$  limit cycle of radius  $\bar{r} = \sqrt{\frac{1}{2}(1 + \sqrt{1+4\alpha})}$   
unique eq. point @ the origin!





$\alpha \rightarrow 0 \rightarrow$  Smaller limit cycles shrink to origin  
 and The bigger one's radius will be  $\frac{1}{2}$ !



Scaling:

EX:

$$\dot{x}_1 = -\alpha x_1 + \beta x_2$$

$$\dot{x}_2 = \frac{\gamma x_1}{\delta + x_1^2} - \eta x_2$$

Greek letters: parameters

Scaled variables:

$$z_i = \frac{x_i}{X_i}, \quad i=1,2 \quad \left\{ \begin{array}{l} T, X_1, X_2 \\ \text{to be determined} \end{array} \right. \quad (\text{constants, "scales"})$$

$$\tau = \frac{t}{T}$$

(2 dependent variables and one independent variable)

$$\frac{dx_i}{dt} = \frac{\partial \tilde{r}}{\partial t} \frac{\partial x_i}{\partial \tilde{r}} = \frac{1}{T} x_i \frac{\partial x_i}{\partial \tilde{r}} = \frac{1}{T} x_i \frac{\partial z_i}{\partial \tilde{r}}$$

$$\rightarrow \frac{dx_i}{dt} = \frac{x_i}{T} \frac{dz_i}{d\tilde{r}} \rightarrow \frac{dz_1}{d\tilde{r}} = \frac{T}{x_1} (-\alpha x_1 z_1 + \beta x_2 z_2)$$

$$\frac{dz_2}{d\tilde{r}} = \frac{T}{x_2} \left( \frac{\gamma x_1 z_1}{\delta + x_1^2 z_1^2} - \eta x_2 z_2 \right)$$

Can Show that:

"with proper choice of  
 $x_1, x_2, T!$ "

$$\begin{cases} \frac{dz_1}{d\tilde{r}} = -a z_1 + z_2 \\ \frac{dz_2}{d\tilde{r}} = \frac{z_1}{1+z_1^2} - b z_2 \end{cases}$$

$\rightarrow$  only two parameters ( $a, b$ ) matter!